# Dynamic Behaviors of $p$-adic Ising-Vannimenus Model on the Cayley Tree of Order Three 

Dogan, M.<br>University of Bahamas, Faculty of Pure and Applied Sciences, School of Mathematics,Physics and Technology, Oekas Field Campus, P.O.Box N-4912, Nassau/Bahamas

E-mail: mutlay.dogan@ub.edu.bs

Received: 12 March 2019
Accepted: 21 March 2020


#### Abstract

Recently, Ising-Vannimenus model on the Cayley tree of order $k=3$ has been studied by Akin (2017) in real case. In this study we continue to investigate Ising-Vannimenus model on the Cayley tree of order $k=3$ in $p$-adic sense. We investigate the dynamic aspects of $p$-adic IsingVannimenus model on the Cayley tree of order $k=3$. We show that the recurrent equation 26 is associated to the model, has four non-trivial fixed points. And one of the fixed points lies in $\mathcal{E}_{p}$ and the rest of fixed points lie in $\mathbb{Z}_{p}^{\star}$. As a main result of the paper, we show that the fixed point $u_{0}$ is attractive and the other fixed points $u_{1}, u_{2}, u_{3}$ are repellent when $u_{i}=p-1$, and neutral when $u_{i} \neq p-1$.


Keywords: $p$-adic Gibbs measures, $p$-adic dynamical systems, IsingVannimenus model and Cayley tree.

## 1. Introduction

It is clear that Ising Vannimenus (see Vannimenus (1981)) model is one of the most crucial models in statistical mechanics. In this work we continue the investigation of Ising-Vannimenus model on the Cayley tree of order $k=3$. In recent studies, existence of $p$-adic quasi Gibbs measures and phase transition were studied in (Mukhamedov et al. (2016), Mukhamedov et al. (2014)) for the $p$-adic Ising-Vannimenus model with competing interactions of nearest, nextnearest and prolonged next-nearest neighbors on the Cayley tree of order $k=2$. In this work, we study the dynamic behaviors of fixed points of $p$-adic IsingVannimenus model with competing interactions of nearest and prolonged nextnearest neighbors on the Cayley tree of order $k=3$.

We consider Hamiltonian ( $H_{n}: \Omega_{V_{n}} \rightarrow \mathbb{Q}_{p}$ ) of the $p$-adic Ising-Vannimenus model as follows;

$$
\begin{equation*}
H(\sigma)=-J \sum_{<x, y>} \sigma(x) \sigma(y)-J_{p} \sum_{>x, y<} \sigma(x) \sigma(y), \tag{1}
\end{equation*}
$$

and $J, J_{p} \in \mathbb{Q}_{p}$ are coupling constants of nearest-neighbor, and prolonged next-nearest-neighbors potentials, respectively. In this paper we cosider $J, J_{p} \neq 0$, i.e. $J \cdot J_{p} \neq 0$. The uniqueness of $p$-adic Gibbs measures in real case was studied in Akin (2017) for the model. In this work, we prove the existence of translation invariant $p$-adic Gibbs measures and dynamic behaviors of fixed points in $p$-adic setting by analyzing the fixed points of dynamical system;

$$
\begin{equation*}
g(u)=\left(\frac{1+c d u}{d+c u}\right)^{3} \tag{2}
\end{equation*}
$$

Note that the results of this paper fails in real setting.

## 2. Preliminaries

## 2.1 -Adic Numbers

In what follows $p$ is a fixed prime number, and $\mathbb{Q}_{p}$ denotes the field of $p$-adic numbers, established by completion of $\mathbb{Q}$ with respect to $p$-adic absolute value $|\cdot|_{p}$. This norm is called non-Archimedean, and satisfies the ultrametric triangle inequality;

$$
\begin{equation*}
|x+y|_{p} \leq \max \left\{|x|_{p},|y|_{p}\right\} . \tag{3}
\end{equation*}
$$

Any nonzero $p$-adic number $x \in \mathbb{Q}_{p}$ can be uniquely shown as

$$
\begin{equation*}
x=p^{\gamma(x)}\left(x_{0}+x_{1} p+x_{2} p^{2}+\ldots\right) \tag{4}
\end{equation*}
$$

where $\gamma=\gamma(x) \in \mathbb{Z}$, and $x_{j}$ are integers such that $0 \leq x_{j} \leq p-1, x_{0}>0$, $j=0,1,2, \ldots$. Here the norm of $x$ is defined by $|x|_{p}=p^{-\gamma(x)}$.

The $p$-adic logarithm is defined by

$$
\begin{equation*}
\log _{p}(x)=\log _{p}(1+(x-1))=\sum_{n=1}^{\infty}(-1)^{n+1} \frac{(x-1)^{n}}{n} \tag{5}
\end{equation*}
$$

and converges when $x \in B(1,1)$,
and $p$-adic exponentials are defined by

$$
\begin{equation*}
\exp _{p}(x)=\sum_{n=1}^{\infty} \frac{x^{n}}{n!} \tag{6}
\end{equation*}
$$

and converges when $x \in B\left(0, p^{-1 /(p-1)}\right)$.
Lemma 2.1. Koblitz (1977). Vladimirov et al. (1994) Let $x \in B\left(0, p^{-1 /(p-1)}\right)$ then we have
$\left|\exp _{p}(x)\right|_{p}=1, \quad\left|\exp _{p}(x)-1\right|_{p}=|x|_{p}<1, \quad\left|\log _{p}(1+x)\right|_{p}=|x|_{p}<p^{-1 /(p-1)}$
and

$$
\log _{p}\left(\exp _{p}(x)\right)=x, \quad \exp _{p}\left(\log _{p}(1+x)\right)=1+x
$$

Lemma 2.2. Khrennikov et al. (2007) If $\left|a_{i}\right|_{p} \leq 1,\left|b_{i}\right|_{p} \leq 1, i=1, \ldots, n$, then

$$
\begin{equation*}
\left|\prod_{i=1}^{n} a_{i}-\prod_{i=1}^{n} b_{i}\right|_{p} \leq \max _{i \leq n}\left\{\left|a_{i}-b_{i}\right|_{p}\right\} \tag{7}
\end{equation*}
$$

We denote the following set,

$$
\begin{equation*}
\mathcal{E}_{p}=\left\{x \in \mathbb{Q}_{p}:|x|_{p}=1, \quad|x-1|_{p}<p^{-1 /(p-1)}\right\} \tag{8}
\end{equation*}
$$

So, from Lemma- 2.1 one concludes that if $x \in \mathcal{E}_{p}$, then there is an element $h \in B\left(0, p^{-1 /(p-1)}\right)$ such that $x=\exp _{p}(h)$. Note that the fundamentals of $p$-adic analysis, $p$-adic mathematical physics were explained in Koblitz (1977), Mahler (1981), Rozikov (2013), Schikhof (1984), Vladimirov et al. (1994).

The $p$-adic integers are defined by

$$
\begin{equation*}
\mathbb{Z}_{p}=\left\{x \in \mathbb{Q}_{p}:|x|_{p} \leq 1\right\}, \tag{9}
\end{equation*}
$$

and the set $\mathbb{Z}_{p}^{\star}=\mathbb{Z}_{p}-p \mathbb{Z}_{p}$ is called $p$-adic units.

### 2.2 Dynamical Systems in $\mathbb{Q}_{p}$

In this subsection we briefly recall some standard terminology of theory of dynamical systems (see Khrennikov and Nilsson (2004)). We define following sets,

$$
\begin{aligned}
& B_{r}(a)=\left\{x \in \mathbb{Q}_{p}:|x-a|_{p}<r\right\}, \quad \bar{B}_{r}(a)=\left\{x \in \mathbb{Q}_{p}:|x-a|_{p} \leq r\right\},(10) \\
& B_{r}(a)=\left\{x \in \mathbb{Q}_{p}: \rho<|x-a|_{p}<s\right\} . S_{r}(a)=\left\{x \in \mathbb{Q}_{p}:|x-a|_{p}=r\right\}(11)
\end{aligned}
$$

for $r, s>0(r<s)$ and $a \in \mathbb{Q}_{p}$. It is clear that $\bar{B}_{r}(a)=B_{r}(a) \cup S_{r}(a)$.
A function $f: B_{r}(a) \rightarrow \mathbb{Q}_{p}$ is said to be analytic if it can be represented by

$$
\begin{equation*}
f(x)=\sum_{n=0}^{\infty} \frac{f^{(n)}(a)}{n!}(x-a)^{n}, \quad f \in \mathbb{Q}_{p} \tag{12}
\end{equation*}
$$

and converges uniformly on the ball $B_{r}(a)$.
Consider a dynamical system $(f, B)$ in $\mathbb{Q}_{p}$, where $f: x \in B \rightarrow f(x) \in B$ is an analytic function and $B=B_{r}(a)$ or $\mathbb{Q}_{p}$ (as given above). We denote $x^{(n)}=f^{n}\left(x^{(0)}\right)$, where $x^{0} \in B$ and $f^{n}(x)=\underbrace{f \circ \cdots \circ f(x)}_{n}$.

If $f\left(x^{(0)}\right)=x^{(0)}$ then $x^{(0)}$ is called a fixed point.
Let $x^{(0)}$ be a fixed point of an analytic function $f(x)$ then,

$$
\lambda=\frac{d}{d x}\left(f\left(x^{(0)}\right)\right),
$$

is a formal derivative of $f$. Then the fixed point $x^{(0)}$ is called attractive when $0 \leq|\lambda|_{p}<1$, indifferent or neutral when $|\lambda|_{p}=1$, and repellent when $|\lambda|_{p}>1$.

## 2.3 p-Adic Measure

Let $(X, \mathcal{B})$ be a measurable space, where $\mathcal{B}$ is an algebra of subsets $X$. A function $\mu: \mathcal{B} \rightarrow \mathbb{Q}_{p}$ is said to be a $p$-adic measure if the following equality holds;

$$
\begin{equation*}
\mu\left(\bigcup_{j=1}^{n} A_{j}\right)=\sum_{j=1}^{n} \mu\left(A_{j}\right) \tag{13}
\end{equation*}
$$

for any $A_{1}, \ldots, A_{n} \subset \mathcal{B}$, and $A_{i} \cap A_{j}=\emptyset(i \neq j)$.
A $p$-adic measure is called a probability measure if $\mu(X)=1$.

### 2.4 Cayley Tree

Let $\Gamma_{+}^{k}=(V, L)$ be a semi-infinite Cayley tree of order $k \geq 1$ with the root $x^{(0)}$ (Each vertex has exactly $k+1$ edges, except $x^{(0)}$ which has $k$ edges). Here $V, L$ are the set of vertices and the set of edges respectively. The vertices $x$ and $y$ are called nearest neighbors if there exists an edge connecting them i.e. $d(x, y)=1$, and they are denoted by $l=\langle x, y>$. Two vertices $x, y \in V$ are called nextnearest neighbors, if $d(x, y)=2$. And the next-nearest-neighbors $x$ and $y$ are called prolonged next-nearest neighbors whenever $x \in W_{n-2}$ and $y \in W_{n}$, and denoted by $>x, y<$, or one-level next-nearest-neighbors, if $x, y \in W_{n}$ for some $n$ and denoted by $>\overline{x, y}<$. Here $W_{n}$ is the nth level of Cayley tree, and we determine following sets;

$$
W_{n}=\left\{x \in V \mid d\left(x, x^{(0)}\right)=n\right\}, \bigcup_{m=1}^{n} W_{m}, L_{n}=\left\{l=<x, y>\in L \mid x, y \in V_{n}\right\} .
$$

The set of direct successors of $x$ is defined by

$$
S(x)=\left\{y \in W_{n+1}: d(x, y)=1, x \in W_{n}\right\} .
$$

Recall that any vertex $x \neq x^{(0)}$ has $k+1$ direct successors except $x^{(0)}$ which has $k$.

## 3. $p$-adic Ising-Vannimenus (IV) model and $p$-adic Gibbs measures

In this section, we consider the $p$-adic Ising-Vannimenus model such that spin values $\sigma(x)$ are from the set $\Phi=\{-1,+1\}$, ( $\Phi$ is called a state space), and these values are assigned to the vertices of Cayley tree $\Gamma^{k}=(V, \Lambda)$. A configuration $\sigma$ on $V$ is defined as a function such that $f: x \in V \rightarrow \sigma(x) \in \Phi$. Using a similar manner one can be defined configurations $\sigma_{n}$ and $\omega$ on $V_{n}$ and $W_{n}$, respectively. The set of all configurations on $V$ (resp. $V_{n}, W_{n}$ ) coincides with $\Omega=\Phi^{V}$ (resp. $\Omega_{V_{n}}=\Phi^{V_{n}}, \Omega_{W_{n}}=\Phi^{W_{n}}$ ). One can see that $\Omega_{V_{n}}=\Omega_{V_{n-1}} \times \Omega_{W_{n}}$, and we define their concatenations as follows;

$$
\left(\sigma_{n-1} \vee \omega\right)(x)= \begin{cases}\sigma_{n-1}(x), & \text { if } x \in V_{n-1}, \\ \omega(x), & \text { if } x \in W_{n},\end{cases}
$$

for configurations $\sigma_{n-1} \in \Omega_{V_{n-1}}$ and $\omega \in \Omega_{W_{n}}$. It is clear that $\sigma_{n-1} \vee \omega \in \Omega_{V_{n}}$.

The Hamiltonian ( $H_{n}: \Omega_{V_{n}} \rightarrow \mathbb{Q}_{p}$ ) of the $p$-adic Ising-Vannimenus model is defined as follows;

$$
\begin{equation*}
H(\sigma)=-J \sum_{<x, y>} \sigma(x) \sigma(y)-J_{p} \sum_{>x, y<} \sigma(x) \sigma(y), \tag{14}
\end{equation*}
$$

with interactions of the nearest-neighbors, and next-nearest-neighbors respectively. Here $J, J_{p} \in \mathbb{Q}_{p}$ are coupling constants.

Let $\mathbf{h}:<x, y>\rightarrow \mathbf{h}_{x y}=\left(h_{x y,++}, h_{x y,+-}, h_{x y,-+}, h_{x y,--}\right) \in \mathbb{Q}_{p}^{4}$ be a vector valued function on each edge, $L$. We consider a $p$-adic probability measure $\mu_{\mathbf{h}}^{(n)}(\sigma)$ on $\Omega_{V_{n}}$ is defined by

$$
\begin{equation*}
\mu_{\mathbf{h}}^{(n)}(\sigma)=\frac{1}{Z_{n}} \exp _{p}\left[-\beta H_{n}(\sigma)+\sum_{x \in W_{n-1}} \sum_{y \in S(x)} \sigma(x) \sigma(y) h_{x y, \sigma(x) \sigma(y)}\right] \tag{15}
\end{equation*}
$$

for an $n \in \mathbb{N}$, and $\beta=\frac{1}{k T}$ (Boltzmann constant, $k$, and temperature, $T$ ). Here, $\sigma_{n}: x \in V_{n} \rightarrow \sigma_{n}(x)$ is a function, and $Z_{n}$ is a partition function as follows;

$$
\begin{equation*}
Z_{n}=\sum_{\sigma_{n} \in \Omega_{V_{n}}} \exp _{p}\left[-\beta H\left(\sigma_{n}\right)+\sum_{x \in W_{n-1}} \sum_{y \in S(x)} \sigma(x) \sigma(y) h_{x y, \sigma(x) \sigma(y)}\right] \tag{16}
\end{equation*}
$$

We consider increasing subsets of the set of states for one dimensional lattices Fannes and Verbeure (1984) as $\mathfrak{G}_{1} \subset \mathfrak{G}_{2} \subset \ldots \subset \mathfrak{G}_{n} \subset \ldots$, where $\mathfrak{G}_{n}$ is the set of states corresponding to non-trivial correlations between $n$-successive lattice points. $\mathfrak{G}_{1}$ is the set of main field states; and $\mathfrak{G}_{2}$ is the set of Bethe-Peierls states, the latter extending to the so-called Bethe lattices. In the probability theory, all these states correspond to so-called Markov chains with memory of length $n$. In this paper, we will discuss a new method of defining Markov chains with memory length of two on a Cayley tree of order three in the $p$-adic sense using the methods described by Akin (2017) (see for details Fannes and Verbeure (1984)).

Let $x \in W_{n-1}$ for some $n \in \mathbb{N}$ and $S(x)=\{y, z, w\}$, where $y, z, w \in W_{n}$ are the direct successors of $x$. Note that $B_{1}(x)=\{x, y, z, w\}$ is a unit semi-ball with a center $x$, where $S(x)=\{y, z, w\}$. We denote the set of all spin configurations on $V_{n}$ by $\Phi^{V_{n}}$ and the set of all configurations on unit semi-ball $B_{1}(x)$ by $\Phi^{B_{1}(x)}$. One can get that the set $\Phi^{B_{1}(x)}$ consists of sixteen configurations;

$$
\Phi^{B_{1}(x)}=\left\{\left(\begin{array}{ccc}
l & k & j  \tag{17}\\
& i &
\end{array}\right): i, j, k, l \in\{-1,+1\}\right\}
$$

Briefly, we do an appropriate definition for the quantities $h\binom{z, y, w}{x}$ as $h_{B_{1}(x)}$.

In this work we consider $\mathbf{h}: V \backslash\left\{x^{(0)}\right\} \times V \backslash\left\{x^{(0)}\right\} \times V \backslash\left\{x^{(0)}\right\} \rightarrow \mathbb{Q}_{p}^{\Phi}$ is a mapping such that

$$
\begin{equation*}
\mathbf{h}:<x, y, z, w>\rightarrow \mathbf{h}_{B_{1}(x)}=\left(h_{B_{1}(x), \sigma(x) \sigma(y) \sigma(z) \sigma(w)}: \sigma(i) \in\{ \pm 1\}\right) \tag{18}
\end{equation*}
$$

where $h_{B_{1}(x), \sigma(x) \sigma(y) \sigma(z) \sigma(w)} \in \mathbb{Q}_{p}, x \in W_{n-1}$ and $y, z, w \in S(x)$. As a result, we use the function $h_{x y z w, \sigma(x) \sigma(y) \sigma(z) \sigma(w)}$ to define the Gibbs measure of any configuration $\left(\begin{array}{cc}\sigma(z) & \sigma(y) \\ & \sigma(w)\end{array}\right)$ that belongs to $\Phi^{B_{1}(x)}$.

In this section, we present the general structure of Gibbs measures with memory length of two on the Cayley tree of order $k=3$. An arbitrary edge $<x^{(0)}, x^{1}>=\ell \in L$ deleted from a Cayley tree $\Gamma_{1}^{k}$ and $\Gamma_{0}^{k}$ splits into two components: semi-infinite Cayley tree $\Gamma_{1}^{k}$ and semi-infinite Cayley tree $\Gamma_{0}^{k}$.

We define the finite-dimensional Gibbs probability distributions on the configuration space $\Omega^{V_{n}}=\left\{\sigma_{n}=\left\{\sigma(x)= \pm 1, x \in V_{n}\right\}\right\}$ as follows;
$\mu_{\mathbf{h}}^{(n)}(\sigma)=\frac{1}{Z_{n}} \exp _{p}\left[-\beta H_{n}(\sigma)+\sum_{x \in W_{n-1}} \sum_{y, z, w \in S(x)} \sigma(x) \sigma(y) \sigma(z) \sigma(w) h_{B_{1}(x), \sigma(x) \sigma(y) \sigma(z) \sigma(w)}\right]$.
with corresponding partition function which is defined by
$Z_{n}=\sum_{\sigma_{n} \in \Omega_{V_{n}}} \exp _{p}\left[-\beta H\left(\sigma_{n}\right)+\sum_{x \in W_{n-1}} \sum_{y, z, w \in S(x)} \sigma(x) \sigma(y) \sigma(z) \sigma(w) h_{B_{1}(x), \sigma(x) \sigma(y) \sigma(z) \sigma(w)}\right]$,
where $\beta=\frac{1}{k T}$. We obtain a new set of $p$-adic Gibbs measures which is different from previous studies Akin (2017), Ganikhodjaev et al. (2011). We consider a construction of an infinite volume distribution with given finite-dimensional distributions. More exactly, we will attempt to find a probability measure $\mu$ on $\Omega$ which is compatible with given measures $\mu_{\mathbf{h}}^{(n)}$, i.e.

$$
\begin{equation*}
\mu\left(\sigma \in \Omega:\left.\sigma\right|_{V_{n}}=\sigma_{n}\right)=\mu_{\mathbf{h}}^{(n)}\left(\sigma_{n}\right), \quad \text { for all } \sigma_{n} \in \Omega_{V_{n}}, n \in \mathbb{N} \tag{20}
\end{equation*}
$$

Kolmogorov consistency condition for $\mu_{\mathbf{h}}^{n}\left(\sigma_{n}\right), n \geq 1$ is defined as follows

$$
\begin{equation*}
\sum_{\omega \in \Omega_{W_{n}}} \mu_{\mathbf{h}}^{(n)}\left(\sigma_{n-1} \vee \omega\right)=\mu_{\mathbf{h}}^{(n-1)}\left(\sigma_{n-1}\right) \tag{21}
\end{equation*}
$$

for any $\sigma_{n-1} \in \Omega_{V_{n-1}}$.
This condition implies the existence of a unique measure $\mu_{\mathbf{h}}$ defined on $\Omega$ with a required condition 20 . Such a measure $\mu_{\mathbf{h}}$ is called a Gibbs measure with memory length of two for the considered model. We define interaction
energy on $V$ with inner configuration $\sigma_{n-1} \in V_{n-1}$ and boundary condition $\eta \in W_{n}$ as follows;

$$
\begin{align*}
H_{n}\left(\sigma_{n-1} \vee \eta\right)= & -J \sum_{<x, y>\in V_{n-1}} \sigma(x) \sigma(y)-J \sum_{x \in W_{n-1}} \sum_{y \in S(x)} \sigma(x) \eta(y) \\
& -J_{p} \sum_{>x, y<\in V_{n-1}} \sigma(x) \sigma(y)-J_{p} \sum_{x \in W_{n-2}} \sum_{z \in S^{2}(x)} \sigma(x) \eta(z) \\
= & H_{n}\left(\sigma_{n-1}\right)-J \sum_{x \in W_{n-1}} \sum_{y \in S(x)} \sigma(x) \eta(y) \\
- & J_{p} \sum_{x \in W_{n-2}} \sum_{z \in S^{2}(x)} \sigma(x) \eta(z) . \tag{22}
\end{align*}
$$

If we follow the same processes as in Akin (2017), then easily we can get the following basic equations. And we express the vector-valued function given in (18) as follows:

$$
\begin{equation*}
\mathbf{h}(x)=\left(h_{1}, h_{2}, h_{3}, h_{4}, h_{5}, h_{6}, h_{7}, h_{8}\right) . \tag{23}
\end{equation*}
$$

Assume that $a=p^{\beta J}$ and $b=p^{\beta J_{p}}, u_{i}^{\prime}=p^{h_{B_{1}(x)}}$ for $x \in W_{n-1}$ and $u_{i}=p^{h_{B_{1}(x)}}$ for $x \in W_{n}$.

For simplicity, when we apply the same technique like in Akin (2017), then we get the following nonlinear dynamical function;

$$
\begin{equation*}
f(x)=\left(\frac{1+(a b)^{2} v_{4}^{4}}{b^{2}+a^{2} v_{5}^{4}}\right)^{3} \tag{24}
\end{equation*}
$$

To reduce (24), we consider that $p^{\frac{2 J}{T}}=a^{2}=c$ and $p^{\frac{2 J_{p}}{T}}=b^{2}=d$, where $T$ is an absolute temperature. Therefore, 24 is conjugate to the following function;

$$
\begin{equation*}
g(x)=\left(\frac{1+c d x}{d+c x}\right)^{3} \tag{25}
\end{equation*}
$$

Hereafter we analyze the equation (25) for the existence of transition invariant $p$-adic Gibbs measures for the considered model.

## 4. Translation-invariant $p$-adic Gibbs measures

In this section, we investigate the existence of translation-invariant $p$-adic Gibbs measures (TIpGM) through analyzing the equation 25. Note that a
function $\mathbf{h}=\left\{h_{B_{1}(x), \sigma_{i}^{(1)}}: i \in\{1,2, \ldots, 16\}\right\}$ is considered as translationinvariant, if $h_{B_{1}(x), \sigma_{i}^{(1)}}=h_{B_{1}(y), \sigma_{i}^{(1)}}$ for all $y \in S(x)$ and $i \in\{1,2, \ldots, 16\}$. A translation-invariant Gibbs measure is defined as a measure $\mu_{\mathrm{h}}$ corresponding to a translation-invariant function $\mathbf{h}$ (see for details Ganikhodjaev et al. (2011), Rozikov (2013)).

After obtaining the equation (25), the existence of $p$-adic Gibbs measures is reduced to the existence of the fixed points of 25 . Therefore, to show the existence of $p$-adic Gibbs measures for all $p>3$, we analyze the following equation (26) which is obtained from (25) with assuming $x=u$,

$$
\begin{equation*}
g(u)=\left(\frac{1+c d u}{d+c u}\right)^{3} \tag{26}
\end{equation*}
$$

where $c, d \in \mathcal{E}_{p}$.
We state the following Proposotion 4.1 to show the existence of $p$-adic Gibbs measures.

Proposition 4.1. Let $p>3$, if $u, v \in \mathcal{E}_{p}$ then following statements hold;
i. $g\left(\mathcal{E}_{p}\right) \subset \mathcal{E}_{p}$, and
ii. $|g(u)-g(v)|_{p} \leq \frac{1}{p}|u-v|_{p}$.

Proof. Let $c, d, u \in \mathcal{E}_{p}$, and using strong triangle inequality, Lemma 2.1 and Lemma 2.2 then easily we get the following items.
i. Let

$$
\begin{aligned}
|g(u)|_{p}=\left|\left(\frac{1+c d u}{d+c u}\right)^{3}\right|_{p} & =\frac{|1+c d u|_{p}^{3}}{|d+c u|_{p}^{3}} \\
& =\frac{|(c d u-1)+2|_{p}^{3}}{|(d-1)+(c u-1)+2|_{p}^{3}}=1 .
\end{aligned}
$$

Then $|g(u)|_{p}=1, u \in \mathcal{E} \mathcal{E}_{p}$.
We show the inequality below,

$$
\begin{aligned}
|g(u)-1|_{p} & =\left|\left(\frac{1+c d u}{d+c u}\right)^{3}-1\right|_{p} \\
& =\left\lvert\, \frac{\left|1+3 c d u+3 c^{2} d^{2} u^{2}+c^{3} d^{3} u^{3}-d^{3}-3 c d^{2} u-3 c^{2} d u^{2}-c^{3} u^{3}\right|_{p}}{|(d-1)+(c u-1)+2|_{p}^{3}} \leq \frac{1}{p}<1 .\right.
\end{aligned}
$$

Hence $g(u) \in \mathcal{E}_{p}$ since (8).
ii. Let $u, v \in \mathcal{E}_{p}$ then the following inequality holds;

$$
\begin{aligned}
|g(u)-g(v)|_{p} & =\left|\left(\frac{1+c d u}{d+c u}\right)^{3}-\left(\frac{1+c d v}{d+c v}\right)^{3}\right|_{p} \\
& =\left|(1+c d u)^{3}(d+c v)^{3}-(1+c d v)^{3}(d+c u)^{3}\right|_{p} \\
& =\left|\left(c d^{2}-c\right)(u-v)\right|_{p} \\
& \leq \frac{1}{p}|u-v|_{p}
\end{aligned}
$$

Hence the proof is completed.

From the Propostion 4.1, the function $g$ satisfies the Banach contraction principle. This means that (26) has a fixed point $u_{0} \in \mathcal{E}$. To find out the other fixed points of (26), let $u=g(u)$ then one gets;

$$
\begin{equation*}
c^{3} u^{4}+\left(3 c^{3} d-c^{3} d^{3}\right) u^{3}+\left(3 c d^{2}-3 c^{2} d^{2}\right) u^{2}+\left(d^{3}-3 c d\right) u-1=0 \tag{27}
\end{equation*}
$$

where $c, d \in \mathcal{E}_{p}$. It is clear that 27) has a solution $u_{0} \in \mathcal{E}_{p}$ since Proposition 4.1 then we rewrite 27) as follows;

$$
\begin{equation*}
\left(u-u_{0}\right)\left[c^{3} u^{3}+\left(3 c^{2} d-c^{3} d^{3}-c^{3} u_{0}\right) u^{2}+A u+B\right]=0 \tag{28}
\end{equation*}
$$

where

$$
A=3 d^{2} c+3 c^{2} d u_{0}-3 c^{2} d^{2}-c^{3} d^{3} u_{0}-c^{3} u_{0}^{2}
$$

and

$$
B=d^{3}+3 d^{2} c u_{0}+3 c^{2} d u_{0}^{2}-3 c d-3 c^{2} d^{2} u_{0}-c^{3} d^{3} u_{0}^{2}-c^{3} u_{0}^{3} .
$$

One of the fixed points of (27) is $u_{0} \in \mathcal{E}_{p}$ and to find the other fixed points of (27) we solve the equation below;

$$
\begin{equation*}
c^{3} u^{3}+\left(3 c^{2} d-c^{3} d^{3}-c^{3} u_{0}\right) u^{2}+A u+B=0, \tag{29}
\end{equation*}
$$

where $c, d, u_{0} \in \mathcal{E}_{p}$. When we divide both sides of (29) by $c^{3}$ then one gets;

$$
\begin{equation*}
u^{3}+\left(\frac{3 c^{2} d-c^{3} d^{3}}{c^{3}}\right) u^{2}+\frac{A}{c^{3}} u+\frac{B}{c^{3}}=0 . \tag{30}
\end{equation*}
$$

In (30) let $a=3 c^{-1} d-d^{3}-u_{0}, b=3 c^{-2} d^{2}+3 c^{-1} d u_{0}-3 c^{-1} d^{2}-d^{3} u_{0}-u_{0}^{2}$ and

$$
e=\frac{d^{3}+3 c d^{2} u_{0}+3 c^{2} d u_{0}^{2}-3 c d-3 c^{2} d^{2} u_{0}-c^{3} d^{3} u_{0}^{2}-u_{0}^{3}}{c^{3}},
$$

then we obtain the following equation;

$$
\begin{equation*}
u^{3}+a u^{2}+b u+e=0 . \tag{31}
\end{equation*}
$$

When we apply the transformation of $u=u-\frac{a}{3}$ then we get the depressed cubic equation;

$$
\begin{equation*}
u^{3}+\left(b-\frac{1}{3} a^{2}\right) u-\left(\frac{2}{27} a^{3}-\frac{1}{3} a b+e\right)=0 . \tag{32}
\end{equation*}
$$

Let $\tilde{a}=b-\frac{1}{3} a^{2}, \tilde{b}=\frac{2}{27} a^{3}-\frac{1}{3} a b+e$ then one gets

$$
\begin{equation*}
u^{3}+\tilde{a} u-\tilde{b}=0 . \tag{33}
\end{equation*}
$$

Hereafter we investigate the existence and number of solutions of depressed equation (33) in $\mathbb{Z}_{p}^{\star}, \mathbb{Z}_{p}, \mathbb{Q}_{p}$ with $p>3$ as in Mukhamedov and Omirov (2014b) and Saburov and Khameini (2015).

Any non-zero $p$-adic number $x \in \mathbb{Q}_{p}$ can be uniquely represented by $x=$ $\frac{x^{\star}}{|x|_{p}}$. Hence, we can represent $\tilde{a}=\frac{\tilde{a}^{\star}}{|\tilde{a}|_{p}}$, and $\tilde{b}=\frac{\tilde{b}^{\star}}{|\tilde{b}|_{p}}$ where $\tilde{a}^{\star}=a_{0}+a_{1} p+$ $a_{2} p^{2}+\ldots, \tilde{b}^{\star}=b_{0}+b_{1} p+b_{2} p^{2}+\ldots$, for any non-zero $\tilde{a}, \tilde{b}$.

Let $D_{0}=-4 a_{0}^{3}-27 b_{0}^{2}$ and $u_{n+3}=b_{0} u_{n}-a_{0} u_{n+1}$ with $u_{1}=0, u_{2}=-a_{0}$, and $u_{3}=b_{0}$ for $n=\overline{1, p-3}$. Accordingly we find the fixed points of (33) through following proposition.

Proposition 4.2. Mukhamedov and Omirov (2014a) Let $p>3$ be a prime. If $|\tilde{a}|_{p}=|\tilde{b}|_{p}=1$ and $|D|_{p}=1$ then (33) has three solutions in $\mathbf{Z}_{\mathbf{p}}^{\star}$, where $\tilde{a}, \tilde{b} \in \mathbf{Q}_{\mathbf{p}}$ with $\tilde{a} \tilde{b} \neq 0$ and $D=-4 \tilde{a}^{3}-27 \tilde{b}^{2}$.

Proof. Using strong triangle inequality and Lemma 2.1
i. Let

$$
\begin{aligned}
|\tilde{a}|_{p} & =\left|b-\frac{1}{3} a^{2}\right|_{p} \\
& =\left|\frac{9 c^{4} d^{2}+9 c^{5} d u_{0}-9 c^{5} d^{2}-3 c^{6} d^{3} u_{0}-3 c^{6} u_{0}^{2}-9 c^{2} d^{2}-c^{6} d^{6}+6 c^{4} d^{4}}{3 c^{6}}\right|_{p} \\
& =\frac{\left|9 c^{4} d^{2}+9 c^{5} d u_{0}-9 c^{5} d^{2}-3 c^{6} d^{3} u_{0}-3 c^{6} u_{0}^{2}-9 c^{2} d^{2}-c^{6} d^{6}+6 c^{4} d^{4}\right|_{p}}{\left|3 c^{6}\right|_{p}} \\
& =1 .
\end{aligned}
$$

ii. Let

$$
\begin{aligned}
|\tilde{b}|_{p}= & \left|\frac{2}{27} a^{3}-\frac{1}{3} a b+e\right|_{p} \\
= & \left\lvert\, \frac{2}{27}\left(\frac{3 c d-c^{3} d^{3}}{c^{3}}\right)\right. \\
& +\frac{d^{3}+3 c d^{2} u_{0}+3 c^{2} d u_{0}^{2}-3 c d-3 c^{2} d^{2} u_{0}-c^{3} d^{3} u_{0}-u_{0}^{3}}{c^{3}} \\
& -\left.\frac{1}{3} \frac{9 c^{2} d^{3}+9 c^{3} d^{2} u_{0}-9 c^{3} d^{3}-3 c^{4} d^{4} u_{0}-3 c^{4} d u_{0}^{2}-3 c^{4} d^{5}}{c^{6}}\right|_{p} \\
& +\left.\frac{1}{3} \frac{-3 c^{5} d^{4} u_{0}+3 c^{5} d^{5}+c^{6} d^{6} u_{0}+c^{6} d^{3} u_{0}^{2}}{c^{6}}\right|_{p} \\
= & \max \left\{\left|\frac{2}{27}\left(\frac{3 c d-c^{3} d^{3}}{c^{3}}\right)\right|_{p}, \mid\right. \\
& \left.\frac{d^{3}+3 c d^{2} u_{0}+3 c^{2} d u_{0}^{2}-3 c d-3 c^{2} d^{2} u_{0}-c^{3} d^{3} u_{0}-u_{0}^{3}}{c^{3}}\right|_{p}, \mid \\
& \frac{1}{3} \frac{9 c^{2} d^{3}+9 c^{3} d^{2} u_{0}-9 c^{3} d^{3}-3 c^{4} d^{4} u_{0}-3 c^{4} d u_{0}^{2}}{c^{6}} \\
& +\left.\frac{1}{3} \frac{-3 c^{4} d^{5}-3 c^{5} d^{4} u_{0}+3 c^{5} d^{5}+c^{6} d^{6} u_{0}+c^{6} d^{3} u_{0}^{2}}{c^{6}}\right|_{p} \\
= & .
\end{aligned}
$$

iii. Let

$$
\begin{aligned}
|D|_{p} & =\left|-4 \tilde{a}^{3}-27 \tilde{b}^{2}\right|_{p} \\
& =\max \left\{\left|-4 \tilde{a}^{3}\right|_{p},\left|27 \tilde{b}^{2}\right|_{p}\right\} \\
& =1
\end{aligned}
$$

From [i], [ii] and [iii] we obtain the required one.

Consequently, the depressed equation (33) has three non-trivial fixed points in $\mathbb{Z}_{p}^{\star}$. So (26) has one fixed point in $\mathcal{E}_{p}$ and three non-trivial fixed points in $\mathbb{Z}_{p}^{\star}$. This result yields the existence of translation invariant $p$-adic quasi Gibbs measures.

## 5. Dynamic Behavior of Fixed Points of Dynamical System

In this section, we investigate the dynamic behavior of the fixed points of (26). In the previous section we proved that the equation (26) has one fixed point in $\mathcal{E}_{p}$ and three fixed points in $\mathbb{Z}_{p}^{\star}$. Here we are going to determine that the fixed points of (26) are attractive, repellent or neutral. To achieve this we state following theorem.

Theorem 5.1. Let $c, d \in \mathcal{E}_{p}$ and $p>3$. The function $f(26)$ has four fixed points $u_{0}, u_{1}, u_{2}, u_{3}$ such that $u_{0}$ is attractive, the fixed points $u_{1}, u_{2}, u_{3}$ are repellent, if $u_{i}=p-1$, and neutral, if $u_{i} \neq p-1$, for $i=1,2,3$, where $u_{0} \in \mathcal{E}_{p}$ and $u_{1}, u_{2}, u_{3} \in \mathbb{Z}_{p}^{\star}$.

Before starting the proof of theorem, we need to state the lemma below.
Lemma 5.2. If $u_{1}, u_{2}, u_{3} \in \mathbb{Z}_{p}^{\star}$ then following statements hold;
i. $\left|1+u_{i}\right|_{p} \leq \frac{1}{p}$ if $u_{i}=p-1$,
ii. $\left|1+u_{i}\right|_{p}=1$ if $u_{i} \neq p-1$,
where $\left|u_{1}\right|_{p}=\left|u_{2}\right|_{p}=\left|u_{3}\right|_{p}=1$ and $i=1,2,3$.

Proof. i. Let $u_{i}=p-1$ then it is clear that

$$
\left|1+u_{i}\right|_{p}=|p|_{p}=\frac{1}{p}
$$

ii. Let $u_{i} \neq p-1$ then it is clear that

$$
\left|1+u_{i}\right|_{p}=1
$$

since $p \nmid\left(1+u_{i}\right)$

Now we are ready to prove Theorem 5.1.

Proof. Since 26) let

$$
g(u)=\left(\frac{1+c d u}{d+c u}\right)^{3}
$$

then we get derivative of $g$ as follows;

$$
g^{\prime}(u)=3 \frac{c d^{2}+2 c^{2} d^{3} u+c^{3} d^{4} u^{2}-c-2 c^{2} d u-c^{3} d^{2} u^{2}}{(d+c u)^{4}}
$$

where $c, d \in \mathcal{E}_{p}$. Thereafter one gets;

$$
\begin{aligned}
\left|g^{\prime}\left(u_{0}\right)\right|_{p} & =\frac{\left|3 c d^{2}+6 c^{2} d^{3} u_{0}+3 c^{3} d^{4} u_{0}^{2}-3 c-6 c^{2} d u_{0}-3 c^{3} d^{2} u_{0}^{2}\right|_{p}}{\left|d+c u_{0}\right|_{p}^{4}} \\
& =\mid 3\left(c d^{2}-1\right)+3+6\left(c^{2} d^{3} u_{0}-1\right)+6+3\left(c^{3} d^{4} u_{0}^{2}-1\right)+3 \\
& -3(c-1)-3-6\left(c^{2} d u_{0}-1\right)-6-3\left(c^{3} d^{2} u_{0}^{2}-1\right)-\left.3\right|_{p} \\
& =|3+6+3-3-6-3|_{p}=0
\end{aligned}
$$

Therefore, $u_{0}$ is an attractive.

Now let us look for the dynamic behaviors of other fixed points $u_{1}, u_{2}, u_{3} \in$ $\mathbb{Z}_{p}^{\star}$. From Lemma 5.2 and ultrametric triangle inequality we get the items below.
i. Let us take $u_{i} \in \mathbb{Z}_{p}^{\star}, i=1,2,3$ and $u_{i}=p-1$ then we get

$$
\begin{aligned}
\left|g^{\prime}\left(u_{i}\right)\right|_{p} & =\frac{\left|3 c d^{2}+6 c^{2} d^{3} u_{i}+3 c^{3} d^{4} u_{i}^{2}-3 c-6 c^{2} d u_{i}-3 c^{3} d^{2} u_{i}^{2}\right|_{p}}{\left|d+c u_{i}\right|_{p}^{4}} \\
& =\frac{\left|\left(1+u_{i}\right)^{2}-\left(1+u_{i}\right)\right|_{p}}{\left|1+u_{i}\right|_{p}^{4}} \\
& =\frac{\left|\left(1+u_{i}\right)\right|_{p}}{\left|1+u_{i}\right|_{p}^{4}} \\
& =\frac{1}{\left|1+u_{i}\right|_{p}^{3}}>1 .
\end{aligned}
$$

Therefore $u_{i}$ are repellent in the considered case $u_{i}=p-1$.
ii. Let us take $u_{i} \in \mathbb{Z}_{p}^{\star}, i=1,2,3$ and $u_{i} \neq p-1$ then we easily get

$$
\begin{aligned}
\left|g^{\prime}\left(u_{i}\right)\right|_{p} & =\frac{\left|3 c d^{2}+6 c^{2} d^{3} u_{i}+3 c^{3} d^{4} u_{i}^{2}-3 c-6 c^{2} d u_{i}-3 c^{3} d^{2} u_{i}^{2}\right|_{p}}{\left|d+c u_{i}\right|_{p}^{4}} \\
& =\frac{\left|\left(1+u_{i}\right)^{2}-\left(1+u_{i}\right)\right|_{p}}{\left|1+u_{i}\right|_{p}^{4}} \\
& =\frac{\left|\left(1+u_{i}\right)\right|_{p}}{\left|1+u_{i}\right|_{p}^{4}}=1
\end{aligned}
$$

Hence $u_{i}$ are neutral in the case of $u_{i} \neq p-1$.

Consequently, we conclude that $u_{0}$ is an attractive, $u_{i}$ are repellent when $u_{i}=p-1$ and $u_{i}$ are neutral when $u_{i} \neq p-1$ for the function (26).

## 6. Conclusion

In the present paper, we obtained the dynamic function (26) as in Akin (2017) and then we proved the existence of the translation invariant $p$-adic Gibbs measures for $p$-adic Ising-Vannimenus model on the Cayley tree of order $k=3$. We found out that one of fixed points of lies in $\mathcal{E}_{p}$ and the other three fixed points lie in $\mathbb{Z}_{p}^{\star}$. As the dynamic behaviours of the model, we proved that the fixed point $u_{0} \in \mathcal{E}_{p}$ is attractive and the other fixed points $u_{i} \in \mathbb{Z}_{p}^{\star}, i=1,2,3$ are repellent when $u_{i}=p-1$, and are neutral when $u_{i} \neq p-1$.

## References

Akin, H. (2017). Phase transition and gibbs measures of vannimenus model on semi-infinite cayley tree of order three. International Journal of Modern Physics B, 31:1-17.

Fannes, M. and Verbeure, A. (1984). On solvable models in classical lattice systems. Commun. Math. Phys., 96:115-124.

Ganikhodjaev, N. N., Akin, H., Uguz, S., and Temir, S. (2011). On extreme gibbs measures of the vannimenus model. Journal of Statistical Mechanics:Theory and Experiment, 10:1742-5468.

Khrennikov, A., Mukhamedov, F., and Mendes, J. (2007). On p-adic gibbs measures of countable state potts model on the cayley tree. Nonlinearity, 20:2923-2937.

Khrennikov, A. and Nilsson, M. (2004). P-adic deterministic and random dynamics. Dordrecht, Netherlands: Springer. https://doi.org/10.1007/978-1-4020-2660-7.

Koblitz, N. (1977). P-adic numbers, p-adic analysis and zeta-function. New York, NY: Springer-Verlag.

Mahler, K. (1981). P-adic numbers and their functions. Cambridge, United Kingdom: Cambridge University Press.

Mukhamedov, F., Dogan, M., and Akin, H. (2014). Phase transition for the $p$-adic ising-vannimenus model on the cayley tree. Journal of Statistical Mechanics: Theory and Experiment, 2014(10):P10031.

Mukhamedov, F. and Omirov, B. (2014a). On cubic equations over $p$-adic fields. Int. Journal of Number Theory, 10:1171-1100.

Mukhamedov, F. and Omirov, B. (2014b). On cubic equations over P-adic fields. International Journal of Number Theory, 10:1171-1100.

Mukhamedov, F., Saburov, M., and Khakimov, O. (2016). Translationinvariant p-adic quasi-gibbs measures for the isingâvannimenus model on a cayley tree. Theoretical and Mathematical Physics, 187(1):583-602. https://doi.org/10.1134/S0040577916040127.

Rozikov, U. (2013). Gibbs Measures on Cayley Trees. Singapore: World Scientific Publishing Co Pte Ltd.

Saburov, M. and Khameini, M. (2015). On description of all translation invariant P-adic gibbs measures for the potts model on the Cayley tree of order three. Mathematical Physics, Analysis and Geometry, 18:18-26.

Schikhof, W. (1984). Ultrametric Calculus: An Introduction to p-Adic Analysis. New York, NY: Cambridge University Press.

Vannimenus, J. (1981). Modulated phase of an ising system with competing interactions on a Cayley tree. Zeitschrift fur Physik B Condensed Matter, 43:141-148.

Vladimirov, V., Volovich, I., and Zelenov, E. (1994). P-adic Analysis and Mathematical Physics. Singapore: World Scientific Publishing Co Pte Ltd.

